

Z-99 PP AND PS POLARITIES IN CORRELATION OF EVENTS AND APPROXIMATING R_{PS} FOR AVO

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Summary

For most interfaces, R_{PP} and R_{PS} have opposite sign, which means an event usually has the same display polarity on PP and PS seismic sections. In processing and interpreting PS sections, we usually rely on correlation of corresponding PS and PP events. Assuming the wrong polarity relationship can lead to a mistie of half a cycle between corresponding PP and PS events and, consequently, over- or underestimated V_P/V_S ratios for the affected intervals. We examine the conditions that determine the display polarities of PP and PS and show that the sign of R_{PS} depends on the sign of ΔU , where $U = \rho\beta^h$ and h is a function of velocities and densities. We also derive a small-angle R_{PS} approximation that is more accurate than the Aki-Richards expression in many situations, not just for small angles.

Introduction

Correlation of PS and PP events enables determination of preliminary V_P/V_S ratios, needed in the PS processing flow. Typically, this correlation relies on synthetic seismograms generated using log data for V_P , V_S and ρ from a nearby well, if available. These synthetics are then used to match events on PP and PS field sections and, from traveltimes ratios, to calculate V_P/V_S ratios. But these efforts can be frustrated if the logs for V_S or ρ are unavailable. If so, ρ may be estimated from V_P using empirical relationships like Gardner's equation (Gardner et al., 1974) or modified versions thereof for specific lithologies (Castagna et al., 1993). If V_S information is unavailable, estimates of interval V_P/V_S ratios have to be used to create PS synthetic stacks (Lawton & Howell, 1992), which may then be fine-tuned to give optimal correlation with the stacked field data (Miller, 1996). In the absence of any of the logs, there may be polarity errors on particular events on the PS synthetics relative to the PP (Figure 1).

The SEG polarity standard (Thigpen et al., 1975) implies, for a minimum-phase wavelet from a compressive source, that a PP reflection from an interface with a *positive* PP reflection coefficient ($R_{PP} > 0$) will begin with a downward (negative) deflection on the recorded seismogram (Sheriff, 2002). Recommended SEG standards for horizontal-component geophones (Landrum et al., 1994) and subsequent proposed standards (Brown et al., 2002) imply that a PS reflection from the same source, when the interface has a *negative* PS reflection coefficient ($R_{PS} < 0$), will also lead to a downward (negative) deflection on the inline geophone on positive-offset traces. Negative-offset traces, which are flipped in preprocessing, originally have the opposite polarity (e.g. Tessmer & Behle, 1988; Brown et al., 2002). Therefore, because for most interfaces, R_{PP} and R_{PS} have opposite sign, an event usually has the same display polarity on PP and PS. In this paper, we study the 'unusual situation', where events on the two sections display opposite apparent polarities, i.e., where $R_{PS}/R_{PP} > 0$.

Approximations to the Zoeppritz equations

In adopting approximations for polarity relations, we do not want to restrict ourselves to low-contrast interfaces, so we assume small angles of incidence, i . We also assume the polarity relationship at small angles to be representative of that of the stacked events in all but the rarest of cases.

Many approximations to the Knott-Zoeppritz equations governing P-SV waves at a welded interface have been derived (mainly for R_{PP} and R_{PS}), some of the earliest being those of Bortfeld (1961), Richards & Frasier (1976) and Aki & Richards (1980), whose notation we essentially follow. We use the symbol r to represent rock parameters generally (e.g. α , β , ρ , σ or μ) and we express an

arbitrary rock parameter and its change over an interface in terms of r (the average of r_1 and r_2) and Δr (the difference $r_2 - r_1$), where 1 and 2 denote media 1 and 2, respectively.

For R_{PP} , many approximations have been published (e.g. Bortfeld, 1961; Richards & Frasier, 1976; Aki & Richards, 1980; Shuey, 1985; Wang, 1999; Ursenbach, 2002) However, for R_{PP} , we use the zero-offset expression as sufficient for characterizing polarity. Some of the published R_{PS} approximations assume small parameter changes (Aki & Richards, 1980; Xu & Bancroft, 1997; Ursenbach, 2002), some, like ours (below) assume small angles (Bortfeld, 1961; Richards & Frasier, 1976; Zaengle & Frasier, 1993; Wang, 1999; Ramos & Castagna, 2001; Carcuz, 2001; Geldart & Sheriff, 2004). It turns out that the differences between the two are not as great as one might think because in the Taylor-series expansions the terms of higher order in i , or $\sin i$, or p (horizontal slowness) tend also to be the terms of higher order in $\Delta r/r$.

Approximating R_{PS}

In deriving an R_{PS} approximation, we start with the exact formula as given by Aki & Richards (1980):

$$R_{PS} = -2 \left[\frac{\cos i_1}{\alpha_1} \left(ab + cd \frac{\cos i_2}{\alpha_2} \frac{\cos j_2}{\beta_2} \right) p \alpha_1 \right] / (\beta_1 D) \quad (1)$$

where a , b , c , d and D are defined by Aki & Richards (1980). These are reprinted by Xu & Bancroft (1997), Ramos & Castagna (2001) and Vant (2003), though Ramos & Castagna have an error in their expression for d . It should be:

$$d = 2(\rho_2 \beta_2^2 - \rho_1 \beta_2^2) \quad \text{and not} \quad d = 2(\rho_2 \beta_2^2 p^2 - \rho_1 \beta_2^2). \quad (2)$$

To get our R_{PS} approximation we first rewrite (1) as:

$$R_{PS} = - \left[\sin 2i_1 \left(ab + cd \frac{\cos i_2}{\alpha_2} \frac{\cos j_2}{\beta_2} \right) \right] / (\alpha_1 \beta_1 D) \quad (3)$$

then apply the small-angle approximation by setting $\sin \theta \approx \theta$ and $\cos \theta \approx 1$ for sines and cosines of incidence angles, i_1 , i_2 , j_1 and j_2 , in the expressions for a , b , c , d and D , giving:

$$R_{PS} = \frac{-\sin 2i_1 (\alpha_2 \beta_2 \rho_2 \Delta \rho + 2\rho_1 \Delta \mu)}{(\rho_1 \alpha_1 + \rho_2 \alpha_2)(\rho_1 \beta_1 + \rho_2 \beta_2)}. \quad (4)$$

If we also approximate the explicit sine factor, it simply reduces from $\sin 2i_1$ to $2i_1$ and we have an expression equivalent to one given by Geldart & Sheriff (2004, p. 70). However this is only a slight simplification and it costs significantly in accuracy at moderate-to-large angles (Figure 2); so we usually choose to retain the explicit sine factor.

Wang (1999) started with the exact formulae for R_{PP} and R_{PS} (Aki & Richards, 1980) and developed Taylor-series expansions in powers of p , and therefore in powers of $\sin \theta$, or θ , i.e., small-angle approximations. His first R_{PS} approximation [his equation (C-3)] is correct up to terms in p^5 . However, to obtain his second, simplified, approximation [his equation (C-5)], he made two assumptions, one of which is quite unjustified [coming from his equation (A-10)]. For R_{PS} this amounts to assuming that $\Delta \rho = 0$, which eliminates one of the two first-order terms in his expressions, even though he retains other terms up to fifth order. Vant (2003) found Wang's first fifth-order approximation [his (C-3)] to be extremely accurate for a selection of interfaces, but the second [his (C-5)] to be quite inaccurate. Truncation of Wang's first R_{PS} approximation after first order gives an expression that is about as accurate as our equation (4) but much more complicated, while the second truncates to an expression that is much less accurate than equation (4).

Conditions for opposite PP and PS polarities

We use the two small-angle approximations to R_{PP} and R_{PS} to determine under what conditions we get opposite display polarities on the same event on PP versus PS data, that is, when $R_{PP}/R_{PS} > 0$. Thus, we start with our own approximation for R_{PS} [equation (4)] and the normal-incidence expression for R_{PP} :

$$R_{PP} = \frac{\rho_2 \alpha_2 - \rho_1 \alpha_1}{\rho_2 \alpha_2 + \rho_1 \alpha_1}. \quad (5)$$

With respect to equation (5):

$$\text{sgn } R_{PP} = \text{sgn}(\rho_2\alpha_2 - \rho_1\alpha_1) = \text{sgn}(\rho\Delta\alpha + \alpha\Delta\rho) = \text{sgn}\left(\frac{\Delta\alpha}{\alpha} + \frac{\Delta\rho}{\rho}\right). \quad (6)$$

With respect to equation (4):

$$\text{sgn } R_{PS} = -\text{sgn}\left\{\alpha_2\beta_2\rho_2 + \rho_1(\beta_2^2 + \beta_1^2)\right\}\Delta\rho + 2\rho_1(\rho_2\beta_2 + \rho_1\beta_1)\Delta\beta = -\text{sgn}\left[f\frac{\Delta\rho}{\rho} + g\frac{\Delta\beta}{\beta}\right] \quad (7)$$

where

$$f = \rho\left[\alpha_2\beta_2\rho_2 + \rho_1(\beta_2^2 + \beta_1^2)\right] \quad \text{and} \quad g = 2\rho_1\beta_1(\rho_2\beta_2 + \rho_1\beta_1). \quad (8)$$

So, for opposite display polarities, or $R_{PP}/R_{PS} > 0$, we need either:

$$\frac{\Delta\alpha}{\alpha} + \frac{\Delta\rho}{\rho} > 0 \quad \text{and} \quad f\frac{\Delta\rho}{\rho} + g\frac{\Delta\beta}{\beta} < 0; \quad \text{i.e.} \quad -\frac{\Delta\alpha}{\alpha} < \frac{\Delta\rho}{\rho} < -\frac{g\Delta\beta}{f\beta} \quad \text{for } R_{PP} > 0 \text{ and } R_{PS} > 0 \quad (9)$$

or:

$$\frac{\Delta\alpha}{\alpha} + \frac{\Delta\rho}{\rho} < 0 \quad \text{and} \quad f\frac{\Delta\rho}{\rho} + g\frac{\Delta\beta}{\beta} > 0; \quad \text{i.e.} \quad -\frac{g\Delta\beta}{f\beta} < \frac{\Delta\rho}{\rho} < -\frac{\Delta\alpha}{\alpha} \quad \text{for } R_{PP} < 0 \text{ and } R_{PS} < 0. \quad (10)$$

Equations (8) to (10) give conditions for occurrence of the unusual situation, that is, reversed display polarity of a reflection event on PS versus PP. We can formulate these conditions from (9) and (10) in another way by first noticing that:

$$\frac{\Delta\rho}{\rho} + \frac{g\Delta\beta}{f\beta} = \frac{\Delta\rho}{\rho} + \Delta\ln\beta^h = \frac{\Delta U}{U}, \quad \text{where} \quad h = \frac{g}{f} \quad \text{and} \quad U = \rho\beta^h. \quad (11)$$

Then we will get the unusual polarity situation, $R_{PP}/R_{PS} > 0$, if either:

$$\frac{\Delta U}{U} < 0 \quad (\text{i.e. } R_{PS} > 0) \quad \text{when} \quad \frac{\Delta Z}{Z} > 0 \quad (\text{i.e. } R_{PP} > 0) \quad (12)$$

or

$$\frac{\Delta U}{U} > 0 \quad (\text{i.e. } R_{PS} < 0) \quad \text{when} \quad \frac{\Delta Z}{Z} < 0 \quad (\text{i.e. } R_{PP} < 0). \quad (13)$$

Conclusions

We have derived a small-angle R_{PS} approximation that is more accurate than the Aki-Richards approximation for three interfaces tested: beyond 25° incidence in two cases and beyond 40° in the third case. We have also derived mathematical expressions for the conditions on the interface rock parameters for an event to be recorded with opposite display polarities on PP and PS data.

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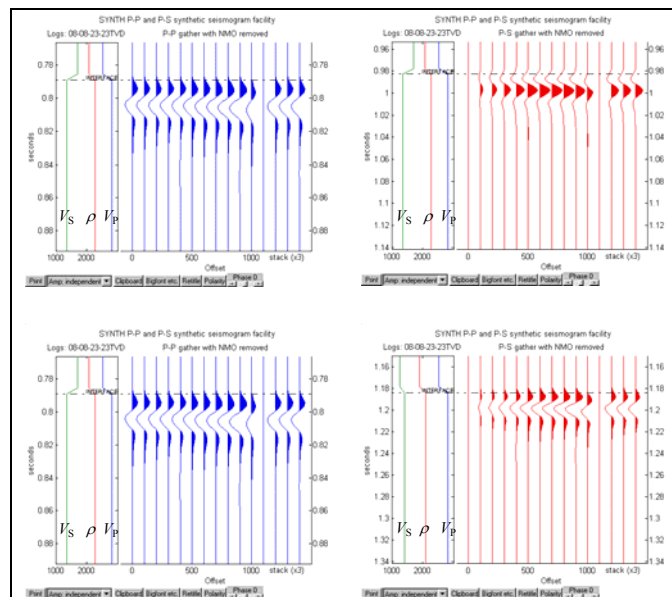


Figure 1. AVO responses and synthetic stacks [PP (blue) and PS (red)] for a model of a sand over a coal, an example of opposite PP-PS display polarities. In modelling this interface (upper figure) all logs were used. A V_p/V_s ratio of 2.0 was used to construct the lower PS synthetics for the scenario of a missing shear-wave log.

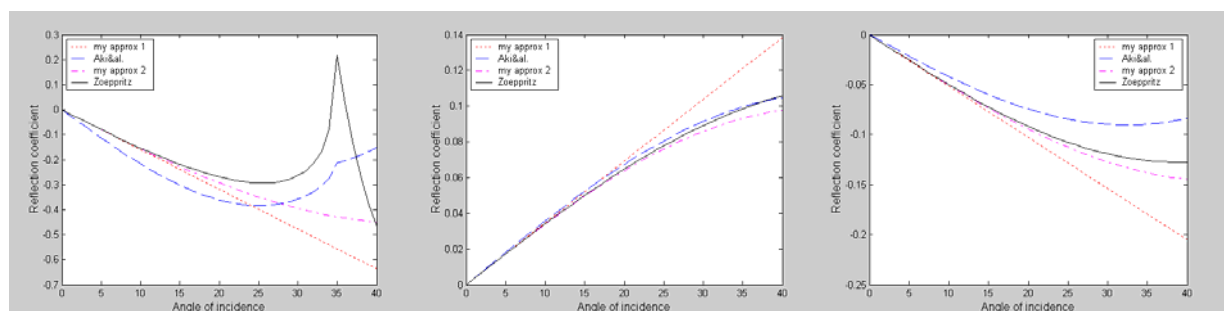


Figure 2. Comparison of the exact R_{PS} curve (solid black) with three approximations: our equation (4) (dashed magenta); Geldart & Sheriff (2004) (dotted red) [equivalent to (4) with $\sin 2i_1 \rightarrow 2i_1$], and Aki & Richards (1980) (coarsely dashed blue). The interface models, left to right, are: young shale over old shale, sandstone over salt, and shale over gas-sand (Brown et al., 2002).