

Efficient Pseudo Gauss-Newton FWI in the time-ray parameter domain

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Introduction

- General Principle of FWI
- * Why FWI fails in industry practice ?

$\boldsymbol{\ast}$ Theory and methods

- Scheme Gradient
- Approximate Hessian
- * Pseudo-Hessian
- Phase Encoded Hessian
- Multiscale Approach
- Numerical Example
- * Conclusions







General Principle of FWI

$$\phi\left(s^{(n)}(\mathbf{r})\right) = \frac{1}{2} \int d\omega \left(\sum_{\mathbf{r}_s, \mathbf{r}_g} \|\delta P\left(\mathbf{r}_g, \mathbf{r}_s, \omega | s^{(n)}(\mathbf{r})\right) \|_2\right)$$

Least-squares misfit function for full waveform inversion







General Principle of FWI

$$s^{(n+1)}(\mathbf{r}) = s^{(n)}(\mathbf{r}) + \mu^{(n)}\delta s^{(n)}(\mathbf{r})$$

The model can updated iteratively

$$\delta s^{(n)} = -\int d\mathbf{r}' H^{(n)-}(\mathbf{r},\mathbf{r}')g^{(n)}(\mathbf{r})$$

The model perturbation can be constructed by gradient and inverse Hessian







Why FWI fails in industry practice?

Extensively computational burden

Slow convergence rate

Cycle skipping problem

Strategies

Source encoding method

Hessian approximations

Multiscale approach







Gradient

$$g^{(n)}(\mathbf{r}) = -\sum_{\mathbf{r}_s, \mathbf{r}_g} \int d\omega \Re \left(\frac{\delta G\left(\mathbf{r}_g, \mathbf{r}_s, \omega | s_0^{(n)}\right)}{\delta s_0^{(n)}(\mathbf{r})} \delta P^* \right)$$

Gradient

$$g^{(n)}(\mathbf{r}) = -\sum_{\mathbf{r}_s, \mathbf{r}_g} \int d\omega \Re \left(\omega^2 \mathcal{F}_s(\omega) G(\mathbf{r}, \mathbf{r}_s, \omega) G(\mathbf{r}_g, \mathbf{r}, \omega) \delta P^* \right)$$

Gradient can be constructed using adjoint state method







Gradient

$$g^{(n)}(\mathbf{r}) = -\sum_{\mathbf{r}_s, \mathbf{r}_g} \int d\omega \Re \left(\frac{\delta G\left(\mathbf{r}_g, \mathbf{r}_s, \omega | s_0^{(n)}\right)}{\delta s_0^{(n)}(\mathbf{r})} \delta P^* \right)$$

Gradient

$$g^{(n)}(\mathbf{r}) = -\sum_{\mathbf{r}_s, \mathbf{r}_g} \int d\omega \Re \left(\omega^2 \mathcal{F}_s(\omega) G(\mathbf{r}, \mathbf{r}_s, \omega) G(\mathbf{r}_g, \mathbf{r}, \omega) \delta P^* \right)$$

RTM image based on crosscorrelation imaging condition







$$g^{(n)}(\mathbf{r}) = -\int d\omega \Re \left\{ \omega^2 \mathcal{F}_s(\omega) \tilde{G}(\mathbf{r}, \mathbf{r}_s, \omega) G(\mathbf{r}_g, \mathbf{r}, \omega) \delta P^* \right\}$$

Linear phase encoded gradient

where $\tilde{G}(\mathbf{r}, \mathbf{r}_s, \omega) = G(\mathbf{r}, \mathbf{r}_s, \omega)e^{i\omega p_s(x'_s - x_s)}$, \mathcal{P}_s is the ray-parameter, x'_s and x_s are the sources' coordinates.



Linear phase encoding strategy

- □ Integration over ray parameter can disperse the crosstalk noise.
- **Different** ray parameters can balance the gradient.







The Functions of Hessian in Least-Squares Inverse Problem



Improve the convergence rate







The Functions of Hessian in Least-Squares Inverse Problem

Hessian Matrix Nonstationary Deconvolution Operator

□ Improve the convergence rate

- **Compensate the geometrical spreading effects and balance the amplitude**
- **Goldson** Suppress the multiple scattering effects and improve the resolution







Approximate Hessian

$$H^{(n)}(\mathbf{r}',\mathbf{r}) = H_1^{(n)} + H_2^{(n)}$$

Full Hessian

$$H_a^{(n)} = H_2^{(n)} \simeq \sum_{\mathbf{r}_s, \mathbf{r}_g} \int d\omega \Re \left\{ \omega^4 G\left(\mathbf{r}_g, \mathbf{r}', \omega | s_0^{(n)}\right) G\left(\mathbf{r}', \mathbf{r}_s, \omega | s_0^{(n)}\right) G^*\left(\mathbf{r}_g, \mathbf{r}'', \omega | s_0^{(n)}\right) G^*\left(\mathbf{r}', \mathbf{r}_s, \omega | s_0^{(n)}\right) \right\}$$

Approximate Hessian in Gauss-Newton Method by Gary and Kris (2011)

$$H_a^{(n)} \simeq diag\left(H_a^{(n)}\right)$$

Diagonal part of the approximate Hessian







Pseudo-Hessian

$$f_{virtual} = -\omega^2 \mathcal{F}_s(\omega) G\left(\mathbf{r}, \mathbf{r}_s, \omega | s_0^{(n)}(\mathbf{r})\right)$$

Virtual Source

$$H_{p_a}^{(n)} = f_{virtual} f_{virtual}^* = \sum_{\mathbf{r}_s} \int d\omega \Re\{\omega^4 | \mathcal{F}_s(\omega)|^2 G(\mathbf{r}', \mathbf{r}_s, \omega) G^*(\mathbf{r}'', \mathbf{r}_s, \omega)\}$$

Pseudo-Hessian

$$I_{dec} = \frac{\sum_{\mathbf{r}_s, \mathbf{r}_g} \int d\omega \Re \left\{ \omega^2 \mathcal{F}_s(\omega) G(\mathbf{r}, \mathbf{r}_s, \omega) G(\mathbf{r}_g, \mathbf{r}, \omega) \delta P^* \left(\mathbf{r}_g, \mathbf{r}_s, \omega \right) \right\}}{\sum_{\mathbf{r}_s} \int d\omega \omega^4 \Re \left\{ |\mathcal{F}_s(\omega)|^2 G\left(\mathbf{r}', \mathbf{r}_s, \omega \right) G^* \left(\mathbf{r}', \mathbf{r}_s, \omega \right) \right\} + \lambda I}$$

Deconvolution imaging condition







Phase Encoded Hessian

$$\begin{aligned} H_{encoded} &= \sum_{\mathbf{r}_s} \int d\omega \Re \left\{ \omega^4 G(\mathbf{r}', \mathbf{r}_s, \omega) G^*(\mathbf{r}'', \mathbf{r}_s, \omega) \right\} \\ &\times \sum_{\mathbf{p}_g} \int d\omega \Re \left\{ G(\mathbf{r}', \mathbf{r}'_g, \omega) e^{i\omega p_g(x'_g - x_{initial})} G^*(\mathbf{r}'', \mathbf{r}_g, \omega) e^{-i\omega p_g(x_g - x_{initial})} \right\} \end{aligned}$$

Receiver-side linear phase encoded Hessian

$$H_{encoded} = H_{exact} + H_{crosstalk}$$

By Tang (2009)

$$H_{encoded} = H_{exact}, \mathbf{p}_g \in (-\infty, +\infty)$$

By Tao and Sen (2013)







Phase Encoded Hessian

$$\begin{split} H_{chirp_encoded} &= \sum_{\mathbf{r}_s} \int d\omega \Re \left\{ \omega^4 G(\mathbf{r}', \mathbf{r}_s, \omega) G^*(\mathbf{r}'', \mathbf{r}_s, \omega) \right\} \\ &\times \sum_{\mathbf{P}_g} \int d\omega \Re \left\{ G(\mathbf{r}', \mathbf{r}'_g, \omega) G^*(\mathbf{r}'', \mathbf{r}_g, \omega) e^{i\omega(p_g + \varepsilon \bigtriangleup p)(x'_g - x_g)} \right\} \end{split}$$

Chirp phase encoded Hessian



Chirp phase encoding strategy





















Hessian Approximations Comparison









Inverse Hessian Comparison









Error Comparison





Gradient Contribution Analysis







Multiscale Approach









Multiscale Approach

- □ Low frequency is responsible to catch the low wavenumber component
- □ High frequency is responsible to add detailed information

$$s\left(\mathbf{r}\right) = s\left(\mathbf{r}\right)^{low} + s\left(\mathbf{r}\right)^{high}$$







Pseudo Gauss-Newton Step

To reduce the computational cost further, we proposed to use one ray parameter in one FWI iteration but change the ray parameter for different iterations.

$$\delta s\left(\mathbf{r}\right) = \frac{\int d\omega \Re \left\{ \omega^{2} \mathcal{F}_{s}(\omega) \tilde{G}(\mathbf{r}, \mathbf{r}_{s}, \omega) G(\mathbf{r}_{g}, \mathbf{r}, \omega) \delta P^{*} \right\}}{diag \left(H_{chirp_encoded} \right) + \lambda I}$$







Pseudo Code for PGN method

BEGIN $\leftarrow s_0$, initial model; **WHILE** $\varepsilon \leq \varepsilon_{min}$ or $n \leq n_{max}$ Identify the ray parameter $p_s^{(n)}$ Identify the frequency band $f^{(n)} = f_0 \rightarrow f_{max}$, $f_{interval}$, every k iterations Generate the data residual δP and apply low-pass filtering $\delta \tilde{P} = \text{low}_\text{pass}(\delta P, f^{(n)})$ Generate the linear phase encoded gradient $g^{(n)}\left(p_s^{(n)}\right)$ **FOR** i = 1 to $\mathbf{p}_s^H, \mathbf{p}_r^H$, every 1 or m iterations Construct the diagonal part of the hybrid phase encoded Hessian $diag\left(H_{en_a}^{(n)}\right)$

END FOR

Calculate the step length $\mu^{(n)}$ using the line search method update the velocity model:

$$s^{(n+1)}(\mathbf{r}) = s^{(n)}(\mathbf{r}) - \mu^{(n)} \left\{ diag \left(H_{en_a}^{(n)} \right) + \lambda I \right\}^{-1} g^{(n)} \left(p_s^{(n)} \right)$$

Calculate the relative least-squares error:

$$\varepsilon = \frac{\|s^{(n)}(\mathbf{r}) - s^{true}(\mathbf{r})\|_2}{\|s^{true}(\mathbf{r})\|_2}$$

END WHILE







Computational Cost Comparison

Table2. Computational cost comparison for different strategies

Methods	$\operatorname{Gradient}$	H_a	$diag(H_{en_a})$	Step length	Cost for one iteration
TGN Method	$2N_s$	$N_s \times N_r$	\	1	$2N_s + N_s \times N_r$
SEGN Method	$2N_p^g$	\setminus	$N_{ps}^H + N_{pr}^H$	1	$2N_p^g + N_{ps}^H + N_{pr}^H$
PGN Method	2	\setminus	$\dot{N_{ps}^H} + \dot{N_{pr}^H}$	1	$N_{ps}^H + \dot{N}_{pr}^H + \dot{2}$





Numerical Experiment





















Effects for Varying Ray Parameter







Sensitivity to the Ray-Parameter Range







Sensitivity to the Encoded Sources









Sensitivity to the Encoded Sources



Initial Velocity Model







Different Hessian Approximations









Different Scaling Methods for FWI













Reverse Time Migration Image Comparison





Image By Initial Velocity Model





Reverse Time Migration Image Comparison





Image By True Velocity Model





Reverse Time Migration Image Comparison





Image By Inverted Velocity Model







Conclusions

- □ Hessian matrix server as a nonstationary deconvolution operator to improve the convergence rate of least-squares inverse problem.
- □ Varying ray-parameter during iterations can reduce the computational cost further and balance the model update.
- □ If the ray-parameter range is too small, the layers with dip angles cannot be inversed in balance, if the ray-parameter range is too large, the convergence rate will be decreased.
- □ If the assembled sources are not dense enough, the crosstalk noise will be very obvious, especially for shallow layers.
- □ Chirp phase encoding strategy can reduce the crosstalk noise better than linear phase encoding strategy with the same number of simulations.
- Diagonal part of the phase encoded Hessian can server as a good approximation of the Hessian to precondition the gradient and increase the convergence rate.
- □ Full waveform inversion with source encoding is efficient for numerical modeling but will increase difficulties in seismic data acquisition and preprocessing.







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