

A least-squares shot-profile application of time-lapse inverse scattering theory

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Goal

Applying least-squares shot-profile migration for imaging time-lapse data sets.

Born approximation of acoustic inverse scattering theory

The forward operator (de-migration) models the scattered seismic wave-field using the Born approximation under the assumption of an acoustic and constant velocity Green's function (Kaplan et al.,2010),

$$\psi_{s}(\mathbf{x}_{g}, z_{g}|\mathbf{x}_{s}, z_{s}; \omega) = f(\omega) \int_{-\infty}^{\infty} G_{0}(\mathbf{x}_{g}, z_{g}|\mathbf{x}', z'; \omega) \left(\frac{\omega}{c_{0}}\right)^{2}$$

$$\alpha(\mathbf{x}', z') G_{0}(\mathbf{x}', z'|\mathbf{x}_{s}, z_{s}; \omega) d\mathbf{x}' dz',$$
(1)

where G_0 is a Green's function for constant acoustic wave-speed c_0 , so,

$$G_0(\mathbf{x}_g, z_g | \mathbf{x}', \mathbf{z}'; \omega) = \frac{1}{2\pi} \int_{-\infty}^{\infty} \left(-\frac{1}{i4k_{gz}} \right) e^{-i\mathbf{k}_{gx} \cdot (\mathbf{x}' - \mathbf{x}_g)} e^{ik_{gz}|z_g - z'|} d\mathbf{k}_{gx}. \tag{2}$$

In equation 1, α is the first order approximation to the scattering potential. Within the context of least-squares migration and SPDR, α is the model (a migrated shot gather). The forward operator in equation 1 describes the mapping from the approximate scattering potential α (the model-space) to the scattered wave-field ψ_s (the data-space) recorded at geophone positions (\mathbf{x}_g, z_g) where $\mathbf{x}_g = (x_g, y_g)$, and due to the seismic source $f(\omega)$ located at (\mathbf{x}_s, z_s) where $\mathbf{x}_s = (x_s, y_s)$. Equation 1 integrates over all possible scattering points (\mathbf{x}', z') where $\mathbf{x}' = (x', y')$. The vertical wave-number k_{gz} in equation 2 is given by the dispersion relation,

$$k_{gz} = sgn(\omega)\sqrt{\frac{\omega^2}{c_0^2} - \mathbf{k}_{gx} \cdot \mathbf{k}_{gx}},$$
 (3)

where $\mathbf{k}_{gx} = (k_{gx}, k_{gy})$ are the lateral wave-numbers (Fourier conjugate variables of $\mathbf{x}_g = (x_g, y_g)$).

Forward (demigration) operator

Let's assume that the earth model is partitioned into n_z layers of constant thickness Δz . First, we define $v_{s(l)}$ for $l = 1 \dots n_z$ such that (Kaplan et al.,2010),

$$V_{s(1)}(\mathbf{k}_{gx}, \omega; \mathbf{x}_s) = u_{p(1)}(\mathbf{k}_{gx}, \omega)g(\mathbf{k}_{gx}, \mathbf{x}_s, \omega)$$

$$V_{s(I)}(\mathbf{k}_{gx}, \omega; \mathbf{x}_s) = \Delta u_p(\mathbf{k}_{gx}, \omega)V_{s(I-1)}(\mathbf{k}_{gx}, \omega; \mathbf{x}_s),$$
(4)

where $u_{p(1)} = \exp(ik_{gz}(z_1 - z_0))/(i4k_{gz})$ and $\Delta u_p = \exp(ik_{gz}\Delta z)$ are phase shift operators and $g(\mathbf{k}_{gx}, \mathbf{x}_s, \omega) = 2\pi f(\omega)e^{-i\mathbf{k}_{gx}\cdot\mathbf{x}_s}$ is the synthetic source term. Second, we define $v_{r(I)}$ so that,

$$V_{r(1)}(\mathbf{k}_{gx},\omega) = u_{p(1)}(\mathbf{k}_{gx},\omega)$$

$$V_{r(I)}(\mathbf{k}_{gx},\omega) = \Delta u_{p}(\mathbf{k}_{gx},\omega)V_{r(I-1)}(\mathbf{k}_{gx},\omega).$$
(5)

Then the forward operator becomes,

$$\psi_{s}(\mathbf{k}_{gx},\omega;\mathbf{x}_{s}) = \left(\frac{\omega}{c_{0}}\right)^{2} \Delta z \sum_{l=1}^{n_{z}} v_{r(l)}(\mathbf{k}_{gx},\omega) \mathcal{F}\left[\mathcal{F}^{*}v_{s(l)}(\mathbf{k}_{gx},\omega;\mathbf{x}_{s})\right] \alpha(\mathbf{x}_{g},z_{l};\mathbf{x}_{s}).$$

The main load of computation is in the two two-dimensional Fourier transforms required per depth and frequency.

Adjoint (migration) operator

The implementation of the adjoint operator is also derived using two iterations. First, we define $v_{s(I)}^*$ for $I = 1 \dots n_z$ so that,

$$\mathbf{v}_{s(1)}^{*}(\mathbf{k}_{gx}, \omega; \mathbf{x}_{s}) = \mathbf{u}_{p(1)}^{*}(\mathbf{k}_{gx}, \omega)\mathbf{g}^{*}(\mathbf{k}_{gx}, \mathbf{x}_{s}, \omega)
\mathbf{v}_{s(I)}^{*}(\mathbf{k}_{gx}, \omega; \mathbf{x}_{s}) = \Delta \mathbf{u}_{p}^{*}(\mathbf{k}_{gx}, \omega)\mathbf{v}_{s(I-1)}^{*}(\mathbf{k}_{gx}, \omega; \mathbf{x}_{s}).$$
(7)

Second, we define $v_{r(I)}^*$ for $I = 1 \dots n_z$ so that,

$$\mathbf{v}_{r(1)}^{*}(\mathbf{k}_{gx}, \omega; \mathbf{x}_{s}) = \mathbf{u}_{p(1)}^{*}(\mathbf{k}_{gx}, \omega)\psi_{s}(\mathbf{k}_{gx}, \omega; \mathbf{x}_{s})$$

$$\mathbf{v}_{r(l)}^{*}(\mathbf{k}_{gx}, \omega; \mathbf{x}_{s}) = \Delta \mathbf{u}_{p}^{*}(\mathbf{k}_{gx}, \omega)\mathbf{v}_{r(l-1)}^{*}(\mathbf{k}_{gx}, \omega; \mathbf{x}_{s}).$$
(8)

Then, the adjoint operator becomes,

$$\alpha^{\dagger}(\mathbf{x}_{g}, z_{l}; \mathbf{x}_{s}) = \Delta\omega \sum_{i} \left(\frac{\omega_{i}}{c_{0}}\right)^{2} \left[\mathcal{F}^{*} \mathbf{v}_{s(l)}^{*}(\mathbf{k}_{gx}, \omega; \mathbf{x}_{s})\right] \mathcal{F}^{*} \mathbf{v}_{r(l)}^{*}(\mathbf{k}_{gx}, \omega; \mathbf{x}_{s}), \quad (9)$$

again requiring two two dimensional Fourier transforms per depth and frequency.

Least-squares time-lapse shot-profile migration

Inversion of difference data: The linear system of equations for the difference data can be expressed as

$$\phi(\mathbf{m}_{diff}) = ||\mathbf{W}_{d}(\mathbf{d}_{diff} - \mathbf{Lm}_{diff})||_{2}^{2} + \mu||\mathbf{W}_{m}\mathbf{m}_{diff}||_{2}^{2}, \tag{10}$$

where \mathbf{m}_{diff} and \mathbf{d}_{diff} are the difference model and data difference sections, respectively.

Joint Inversion of base and monitor data: Let's \mathbf{d}_b and \mathbf{d}_m represent the data sets for base and monitor surveys, respectively. Similarly, \mathbf{m}_b and \mathbf{m}_m represent the subsurface scattering potential at base and monitor survey, respectively. A regularized solution can be found by minimizing the following cost function (Ayeni and Biondi, 2010)

$$J = \left\| \begin{bmatrix} \mathbf{L}_{b} \ \mathbf{0} \\ \mathbf{0} \ \mathbf{L}_{m} \end{bmatrix} \begin{bmatrix} \mathbf{m}_{b} \\ \mathbf{m}_{m} \end{bmatrix} - \begin{bmatrix} \mathbf{d}_{b} \\ \mathbf{d}_{m} \end{bmatrix} \right\|^{2} + \left\| \begin{bmatrix} \mu_{b} \mathbf{R}_{b} \ \mathbf{0} \\ \mathbf{0} \ \mu_{m} \mathbf{R}_{m} \end{bmatrix} \begin{bmatrix} \mathbf{m}_{b} \\ \mathbf{m}_{m} \end{bmatrix} \right\|^{2} + \left\| \begin{bmatrix} \lambda_{b} \mathbf{D}_{b} \lambda_{m} \mathbf{D}_{m} \end{bmatrix} \begin{bmatrix} \mathbf{m}_{b} \\ \mathbf{m}_{m} \end{bmatrix} \right\|^{2},$$

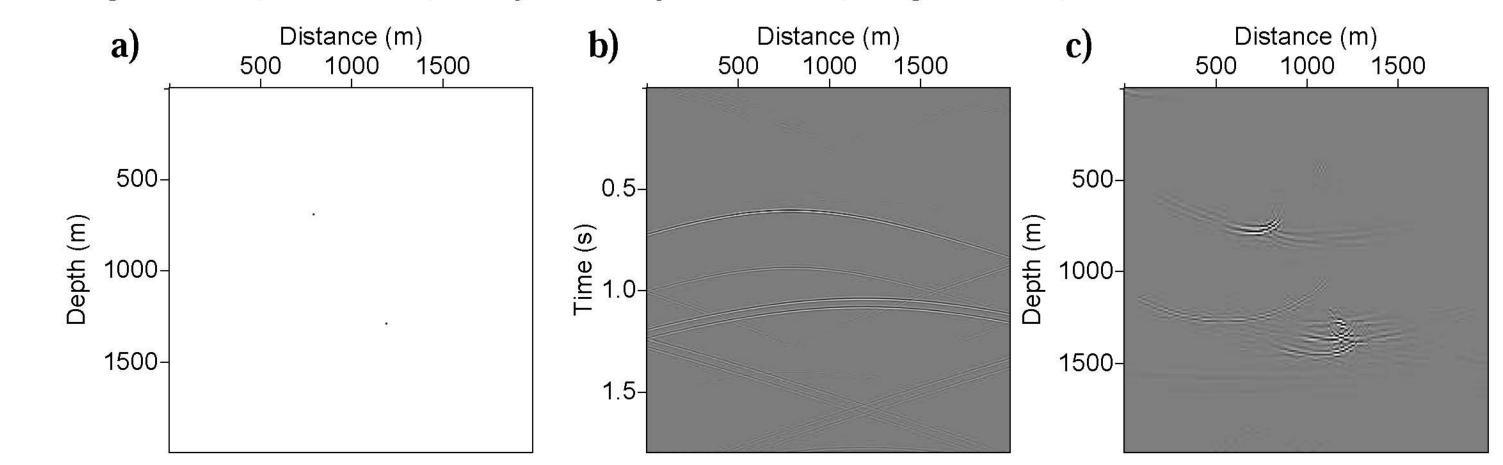
$$+ \left\| \begin{bmatrix} \lambda_{b} \mathbf{D}_{b} \lambda_{m} \mathbf{D}_{m} \end{bmatrix} \begin{bmatrix} \mathbf{m}_{b} \\ \mathbf{m}_{m} \end{bmatrix} \right\|^{2},$$

$$(11)$$

Where \mathbf{R}_b and \mathbf{R}_m are model regularization terms and \mathbf{D}_b and \mathbf{D}_m are temporal regularization terms. The model regularization can be smoothness or sparseness for spatial dimensions and the temporal regularization can be a derivative operator between the surveys.

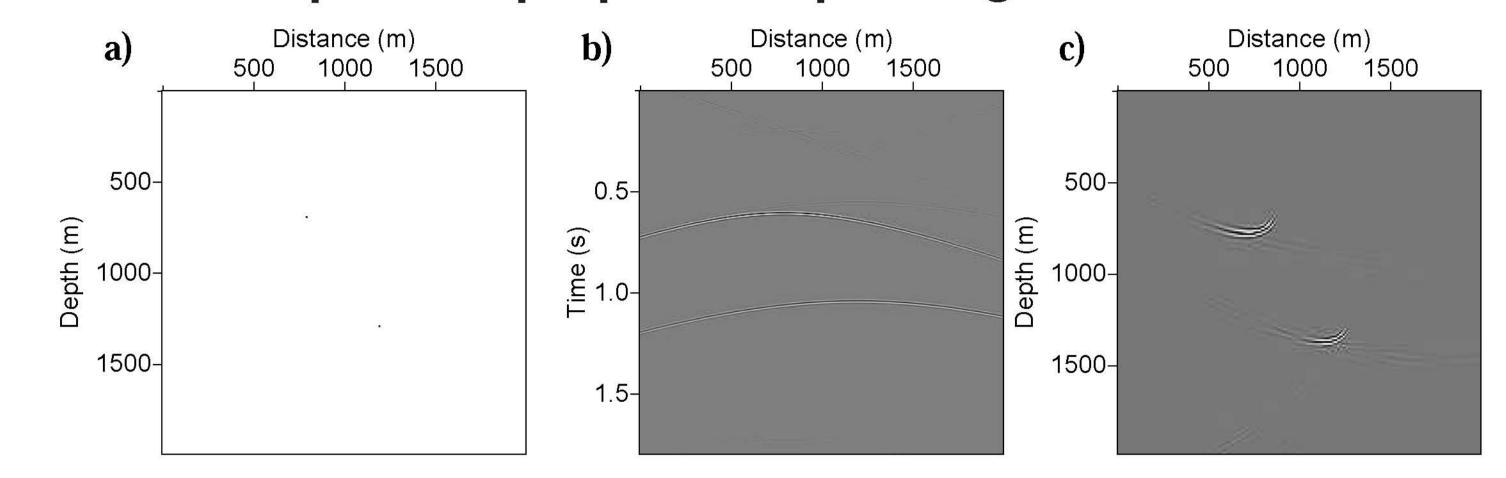
Examples

No zero-padding: a) Scattering potential, b) forward operator (de-migration), and c) adjoint operator (migration).

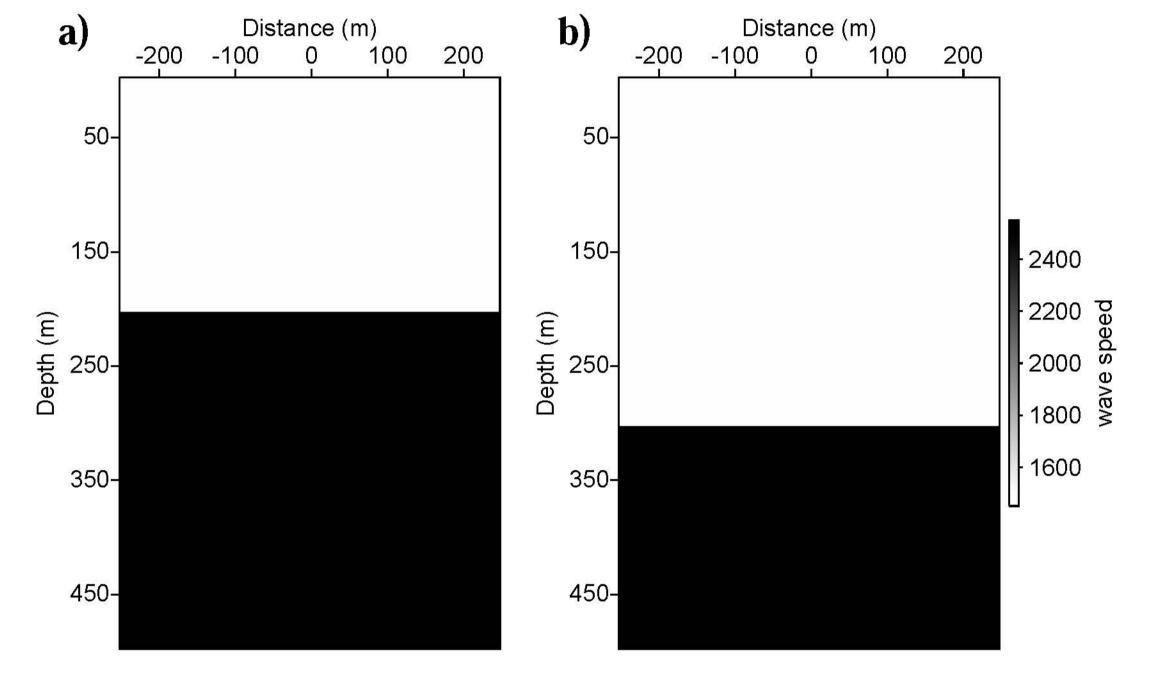


Examples

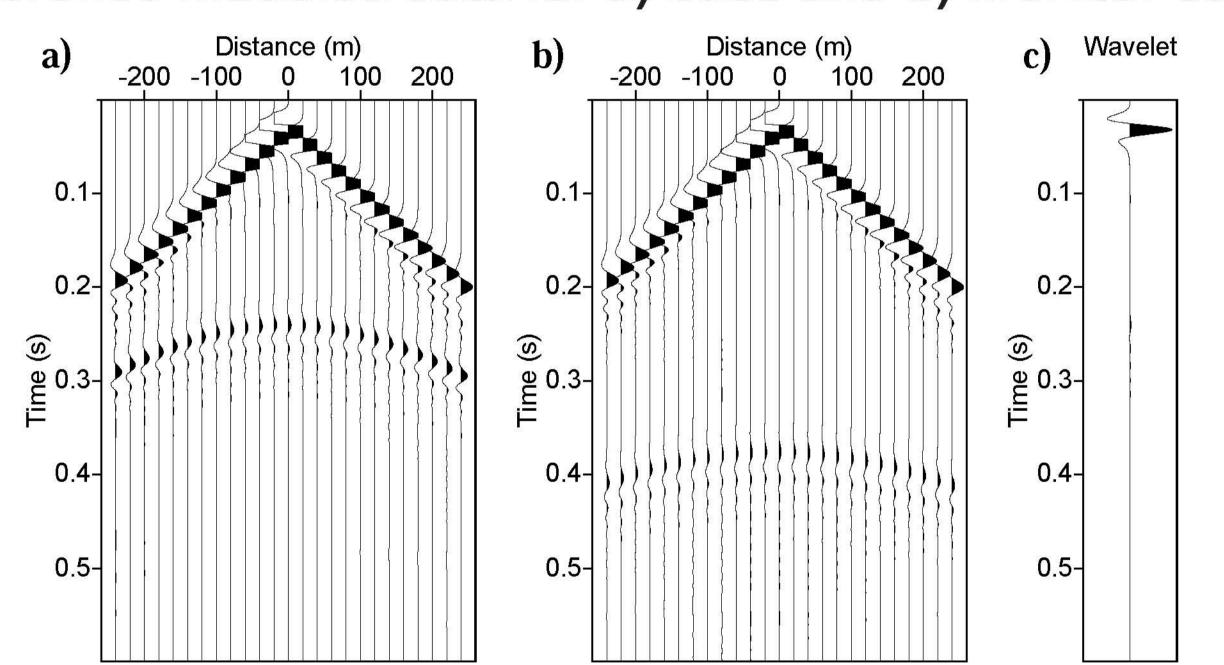
Previous example with proper zero-padding.



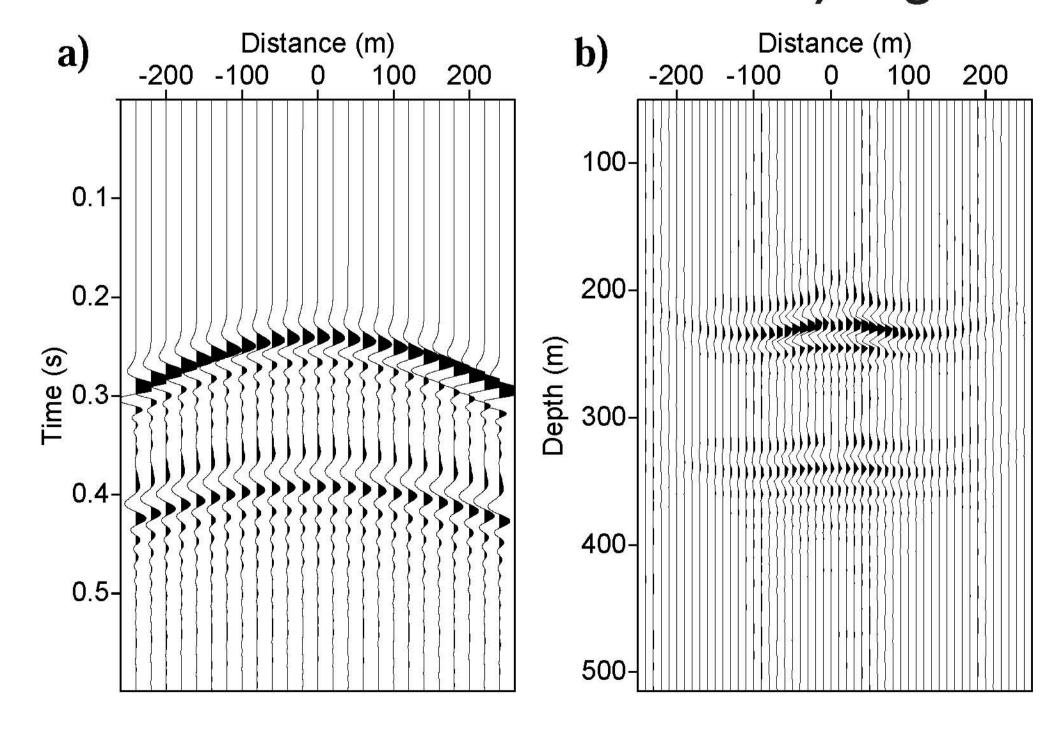
A synthetic velocity model for a) base and b) monitor surveys.



Finite-difference modeled data for a) base and b) monitor surveys.



▶a) The base and monitor difference section. b) Migrated image.



Acknowledgement

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